

Analysis and Design of Filament Power Supply with Voltage-Fed Single-Switch ZVS Inverter

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Abstract. In this paper, an off-line voltage-fed single-switch zero-voltage-switching (ZVS) inverter is analyzed and designed for the filament power supply of cooker magnetrons. Due to the low resistance of cathode filament, an isolation transformer with large turn ratio is used. To obtain the ZVS operation, a parallel-resonant capacitor is connected to the primary of transformer. The capacitor, transformer and filament constitute a parallel resonant tank. Both the transformer leakage inductance and the transistor output capacitance are merged into the resonant tank. The resonant energy provides the ZVS operation of the inverter switch, and the secondary voltage drives the cathode filament. To verify the analysis, a prototype inverter is built and tested. Experimental results validate the circuit design.

1 Introduction

Microwave heating is an important contactless heating approach and the related techniques and applications in household, industrial, agricultural and medical fields are continuously developed [1-3]. The microwave generator of microwave ovens is a continuous-wave magnetron. Its basic components include a hollow metal cylinder, a cathode filament, two ring magnets, even-number anode resonant cavities and an output antenna. The cathode filament is located in the central of the hollow cylinder and is surrounded by the even-number anode cavities. The space between the anode and cathode is the working space of magnetron. In the space, the dc voltage between anode and cathode forms a horizontal electric field and the two ring magnets respectively located on the top and bottom of the hollow cylinder build a vertical magnetic field. As the cathode filament is heated to a critical temperature, the cathode filament emits thermal electrons to the working space. Then, the perpendicular electric and magnetic fields respectively exert forces to the space electrons. As the space electrons sweep across the opening of anode cavities, the polarities of adjacent anode vanes are alternately changed and hence the LC resonance of the anode cavities is caused. At proper conditions, the space electron clouds eventually form rotating spokes. As the spoke number is a half of the anode cavity number, the magnetron enters the pi mode operation. To easily enter the pi mode, the even and the odd anode vanes are respectively tied together. The anode output antenna intercepts the energy of anode resonant cavities and radiates microwaves. With the constant magnetic and electric fields, the anode power is determined by the average anode current.

The cathode filament is the heart of magnetrons. To prolong the usable life, the filament should be preheated

to the thermionic emission temperature before applying negative high voltage to the cathode filament. In general, the pre-heating of cathode filament may take more than two seconds. Inappropriate start-up and filament temperature usually shorten the filament life and degrade the performance. In order to keep filament at proper temperature, the filament driving voltage should be decreased with the average anode current. Shown in Fig.1 is an example that the filament driving voltage decreases with the average anode current. Accordingly, as shown in Fig.2, besides the anode power supply (APS), a complete magnetron driver needs a separate filament power supply (FPS) to preheat and regulate the filament power. In this work, the voltage source V_d is rectified from the utility line. The anode of magnetron is grounded to the heat sink. The APS supplies a negative dc voltage to the cathode and the FPS heats the filament with a small ac voltage. To preheat the filament, switch S1 is turned on before switch S2.

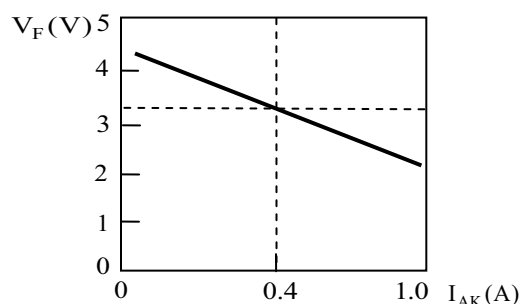


Fig. 1. The required filament voltage decreased with the average anode current.

Currently, there are many types of soft-switching circuits can be used to improve the circuit efficiency and

reduce the volume and weight of power supplies. Due to simplicity and reliability, single-switch zero-voltage-switching (ZVS) circuits have been widely used in induction heating, cooking, LCD TV, EV chargers and other industrial applications [4-10].

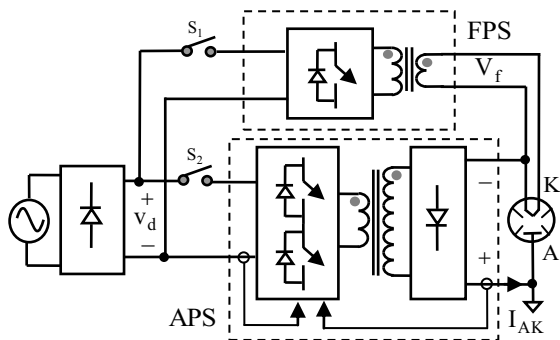


Fig. 2. The block diagram of an off-line magnetron driver.

2 Circuit operations

Fig.3 shows the circuit diagram of the FPS designed in this paper. It is a voltage-fed single-switch ZVS inverter. A step-down transformer is used to attain the required voltage conversion ratio and the secondary voltage drives the cathode filament.

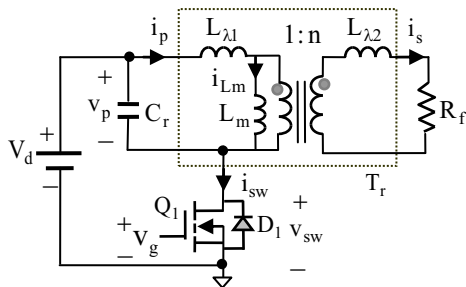


Fig. 3. The circuit diagram of single-switch ZVS inverter.

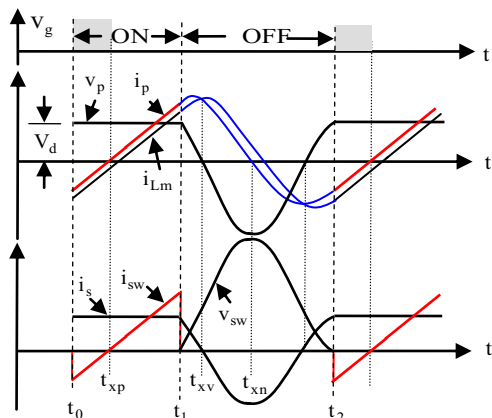


Fig. 4. The principal waveforms of single-switch ZVS inverter.

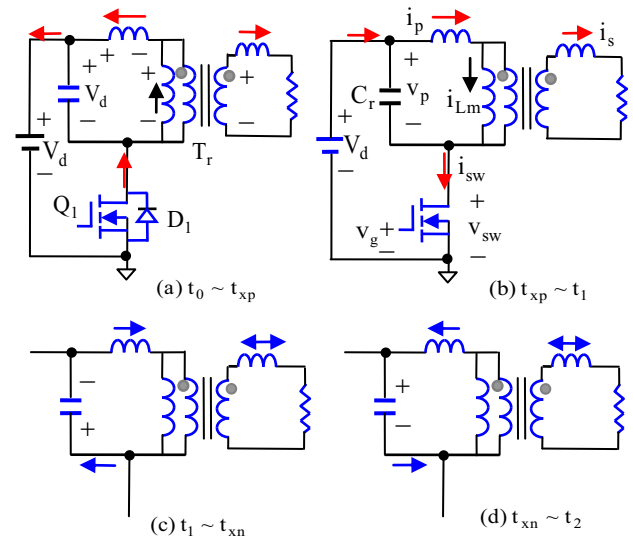


Fig. 5. The circuit operations of the single-switch ZVS inverter.

As switch Q_1 is turned off, transformer T_r and capacitor C_r comprise a parallel resonant tank. The primary leakage inductance and the transistor output capacitance are included in the resonant tank. By low-damped resonance, the resonant energy makes the resonant voltage reach the dc voltage V_d . Thus, switch Q_1 can be turned on at zero-voltage condition and hence switch loss is reduced. Fig.4 conceptually shows the current and voltage waveforms of the ZVS inverter. According to the primary current waveform, a switching cycle is divided into four stages. The detailed circuit operations are sketched in Fig.5.

At ON stage, $t_0 < t < t_1$, the primary voltage equals to V_d . The switch current i_{sw} is negative at t_0 . Since D_1 is conducted, Q_1 is turned on at zero voltage condition. After t_{xp} , switch current rises from zero. As Q_1 is cut off at t_1 , the OFF stage begins. At OFF stage, $t_1 < t < t_2$, T_r resonates with capacitor C_r , and the filament R_f consumes the resonant energy. At t_{xn} , the sinusoidal resonant current crosses zero and the resonant voltage reaches maximum. By small damping ratio, the primary resonant voltage rises to V_d at t_2 . Then, D_1 naturally conducts and Q_1 can be turned on at zero voltage condition.

3 Circuit analyses

Fig.6 shows a simplified resonant circuit, where L_p is the primary inductance, $L_p = L_m + L_{\lambda 1}$, L_m is the magnetizing inductance, $L_m = kL_p$, k is the transformer coupling factor, $k = L_m / (L_m + L_{\lambda 1})$. C_{sw} is the switch output capacitance and C_p is the primary capacitance, $C_p = C_r + C_{sw}$. As the secondary leakage inductance is

ignored, the primary resistance can be approximated as $R_p \approx R_f/n^2$. For driving low resistance filament, the transformer turn ratio is much less than unity, $n = N_s/N_p \ll 1$. The initial inductor current and initial capacitor voltage are I_0 and V_0 respectively.

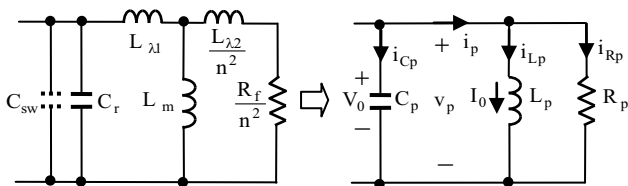


Fig. 6. The simplified equivalent circuit.

With coupling factor k and turn ratio n , the secondary voltage of ON stage can be approximately expressed as

$$v_s(t-t_0) \approx knv_p(t-t_0) = knV_d. \quad (1)$$

The transformer secondary current is

$$i_s(t-t_0) \approx \frac{knV_d}{R_f} + [i_s(t_0) - \frac{knV_d}{R_f}] \exp[-\frac{R_f}{L_{\lambda 2}}(t-t_0)] \quad (2)$$

At OFF stage, the characteristic equation of resonant tank is

$$s^2 + \frac{1}{R_p C_p} s + \frac{1}{L_p C_p} = s^2 + 2\alpha s + \omega_0^2 = 0, \quad (3)$$

Here α is the damping factor, $\alpha = \frac{1}{2R_p C_p}$, ω_0 is the

natural frequency, $\omega_0 = \frac{1}{\sqrt{L_p C_p}}$. The damping ratio is

$\xi = \frac{\alpha}{\omega_0}$. With the loaded quality factor Q_L , (3) can be

rewritten as

$$s^2 + 2\alpha s + \omega_0^2 = s^2 + \frac{\omega_0}{Q_L} s + \omega_0^2 = 0. \quad (4)$$

$$Q_L = \frac{\omega_0}{2\alpha} = \frac{1}{2\xi} = \frac{R_p}{Z_0}, \quad Z_0 = \omega_0 L_p = \frac{1}{\omega_0 C_p} = \sqrt{\frac{L_p}{C_p}}.$$

For under-damped case, $\xi < 1$ or $Q_L > 0.5$, the characteristic equation has two conjugate complex roots,

$$s_{1,2} = -\alpha \pm j\omega_d, \quad \omega_d = \sqrt{\omega_0^2 - \alpha^2}. \quad (5)$$

The natural voltage response of the primary is

$$v_p(\tau) = e^{-\alpha\tau} [B_1 \cos(\omega_d\tau) + B_2 \sin(\omega_d\tau)], \quad \tau = t - t_1. \quad (6)$$

With initial conditions,

$$v_p(0) = V_0, \quad \frac{dv_p(0)}{dt} = \frac{V_0}{R_p C_p} + \frac{I_0}{C_p},$$

the constants B_1 and B_2 are determined.

$$B_1 = V_0, \quad B_2 = V_0 \frac{\alpha}{\omega_d} - (I_0 + \frac{V_0}{R}) \frac{1}{\omega_d C_p}. \quad (7)$$

With the voltage response, the resistive current is

$$i_{Rp}(\tau) = v_p(\tau) / R_p. \quad (8)$$

The inductive current obtained by the KCL is

$$i_{Lp}(\tau) = -i_{Rp}(\tau) - i_{Cp}(\tau), \quad i_{Cp}(\tau) = C_p \frac{dv_p(\tau)}{dt}. \quad (9)$$

For un-damped case, $\alpha = 0$, $B_1 = V_0$, $B_2 = I_0 Z_0$, the primary voltage and current are respectively

$$v_p(\tau) = V_0 \cos(\omega_0\tau) - I_0 Z_0 \sin(\omega_0\tau). \quad (10)$$

$$i_p(\tau) = I_0 \cos(\omega_0\tau) + \frac{V_0}{Z_0} \sin(\omega_0\tau) \quad (11)$$

The primary voltage equation (10) can be written as

$$v_p(\tau) = -V_{pk} \sin(\omega_0\tau - \phi) \quad (12)$$

$$V_{pk} = V_0 \sqrt{1 + \gamma^2}, \quad \gamma = \frac{I_0 Z_0}{V_0}, \quad \phi = \sin^{-1} \left(\frac{1}{\sqrt{1 + \gamma^2}} \right).$$

As shown in Fig.4, the parallel resonance begins at t_1 and ends at t_2 . The resonant voltage crosses zero at t_{xv} and the phase angle between t_{xv} and t_1 is ϕ . The total angle of the OFF stage is $\pi + 2\phi$. As the angle of ON stage is expressed as $\delta\pi$, $0 < \delta < 1$. The relation between δ and γ can be found by the volt-second balance.

$$V_d \delta\pi + 2 \int_0^\phi v_p(\theta - \phi) d\theta = \int_\phi^{\pi+\phi} v_p(\theta - \phi) d\theta - \pi V_d, \quad (13)$$

$\theta = \omega_0 t$. For undamped case, $\alpha = 0$,

$$\delta = \gamma \frac{2}{\pi} - 1. \quad (14)$$

The voltage stress applied on the switch Q_1 is

$$V_{sw} = V_d + V_{pk} = V_d (1 + \sqrt{1 + \gamma^2}) \quad (15)$$

In case that $\delta = 0.3$, $\gamma = 2.04$, $V_{sw} = 3.27 V_d$.

In order to reduce the switch voltage stress, the interval of the ON stage should be decreased and hence a smaller inductance and a larger capacitance are adopted at the desired natural angular frequency.

To easily obtain the ZVS condition, the damping ratio of the RLC tank should be smaller than 0.1. Namely, the loaded quality factor should be greater than 5. In case of $\xi \leq 0.1$, the primary minimum resistance is $R_p \geq 5Z_0$. With the minimum resistance, the condition of the characteristic impedance can be expressed as

$$\sqrt{\frac{L_p}{C_p}} \leq \frac{R_{f,\min}}{5n^2} \quad (16)$$

As the exponential term of (2) is ignored, the energy dissipated at the ON stage can be expressed as

$$W_{ON} = \int_{t_0}^{t_1} i_s^2(t-t_0)R_f dt, \quad i_s(t-t_0) \approx \frac{knV_d}{R_f} \quad (17)$$

At OFF stage, the dissipated energy can be expressed as

$$W_{OFF} = \int_{t_1}^{t_2} i_s^2(\tau)R_f dt, \quad i_s(\tau) \approx \frac{knv_p(\tau)}{R_f}, \quad \tau = t-t_1 \quad (18)$$

4 Simulation and experiment

An off-line single-switch ZVS inverter circuit is designed for a 33W/3V filament. The widely used 8-pin control chip L6561 is used to control the switch conduction time and detect the zero point of switch voltage. The switching frequency is about 20 kHz at rated power and input voltage. To easily obtain the ZVS condition, the loaded quality factor of the RLC parallel resonant tank is set larger than 5. The designed circuit parameters are listed as follows.

$$P_{avg} = 33W \quad , \quad R_{f,\min} = 0.16\Omega \quad , \quad R_f = 0.32\Omega \quad , \\ V_d = 300V \quad , \quad n = 0.02 \quad , \quad k = 0.92 \quad , \quad L_p = 480\mu H \quad , \\ C_p = 0.1\mu F \quad , \quad Z_0 = 69\Omega \quad , \quad f_{sw} \geq 20kHz \quad .$$

Fig.7 displays the natural responses of the primary voltage with initial voltage 300V and initial current 6A. The four traces respectively correspond to the 0.1, 0.16, 0.32 and 100 Ohms filaments. It illustrates that the RLC parallel circuit with above initial conditions can obtain the ZVS operation while the filament is greater than 0.1 Ohm, and the corresponding switch voltage stress at 100Ohms load is about 800V.

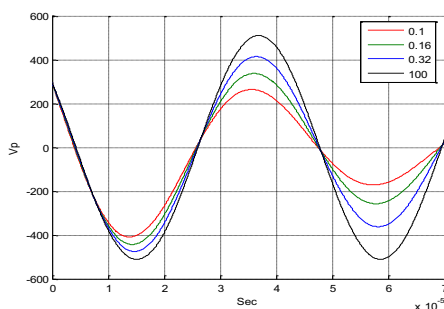


Fig. 7. The simulation of characteristic equation at four different loads.

Fig.8, Fig.9, and Fig.10 respectively show the waveforms simulated at 100 Ω , 0.32 Ω , and 0.16 Ω loads. It can be found that the reflected current increases with the decrease of the filament resistance. As shown in Fig.8, the transformer magnetizing current is equal to the primary current, and the negative switch current that flows through the anti-parallel diode is greater than that of Fig.9 and Fig.10. The switch voltage stress is about 800V. The simulated waveforms demonstrate the circuit analyses. The switch current is a part of the primary current, and the difference between the primary and the magnetizing currents is the reflected load current. The transistor switch obtains the ZVS operation and the sinusoidal resonant waveforms can be observed at the OFF stage. Fig.11 and Fig.12 display the practical waveforms of a scale-down inverter respectively measured at light-load and heavy-load conditions. The ac output voltage of the inverter is rectified and filtered to a dc signal and then fed back to the control chip L6561 to control the conduction time of switch. The output power of the ZVS inverter is dominated by the switch conduction time. And, the switch off time is determined by the resonant parameters. The measured waveforms coincide with the simulated waveforms. The waveforms shown in Fig.13 depict the relation between the filament current and the anode to cathode voltage. At start-up stage, the inverter heats the filament, and the preheating time lasts about 3 seconds. The anode to cathode voltage reaches about 7kV. In order to protect the magnetron from high voltage stress, the anode power supply should be turned on after the pre-heating process.

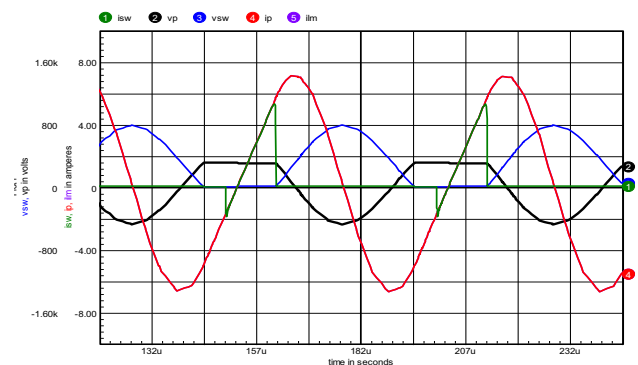


Fig. 8. The waveforms simulated with 100 Ohms load.

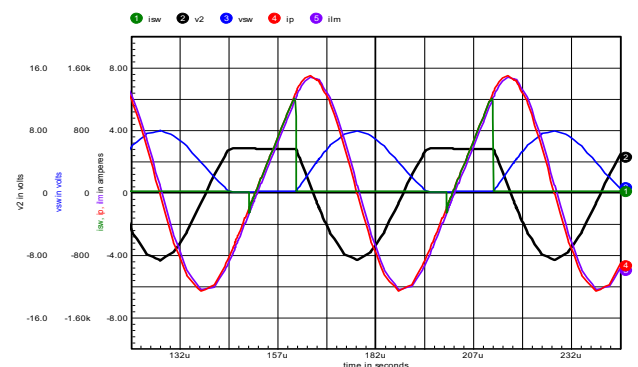


Fig. 9. The waveforms simulated with 0.32 Ohms load.

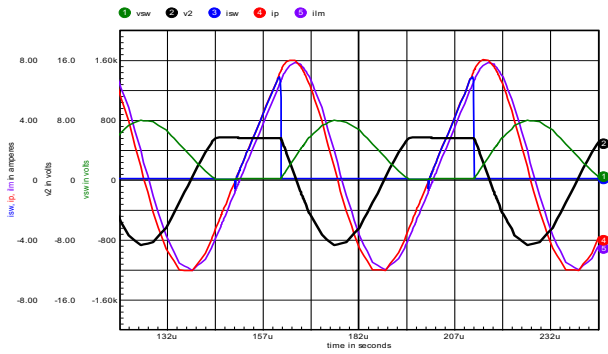


Fig. 10. The waveforms simulated with 0.16 Ohms load.

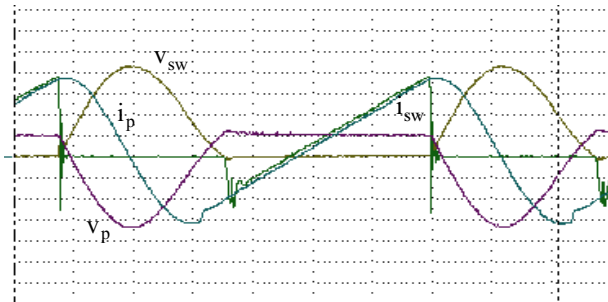


Fig. 11. The steady-state waveforms at light load conditions.

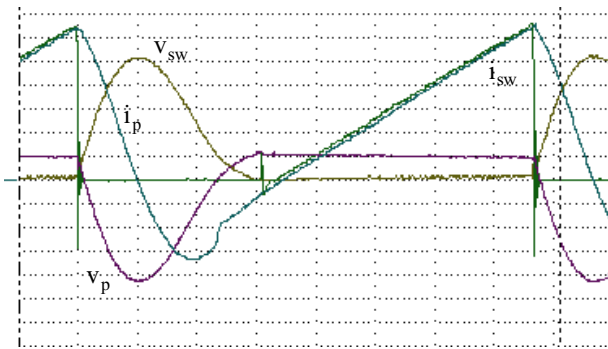


Fig. 12 The steady-state waveforms at heavy load conditions.

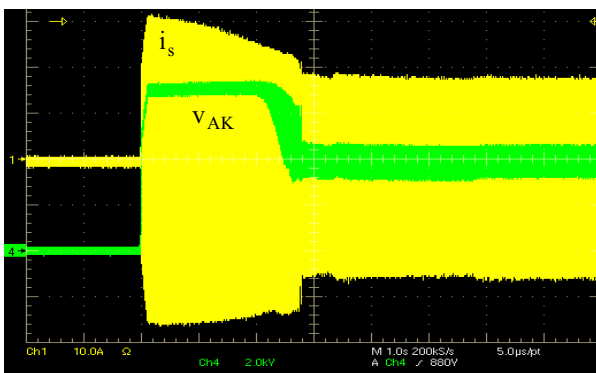


Fig. 13 The waveforms of filament current (10A/Div) and anode to cathode voltage (2kV/Div) at start-up stage, Time:1s/Div.

5 Conclusions

Before applying dc high voltage to a cooker magnetron, the cathode filament is preheated to reach the thermionic emission temperature. As the magnetron starts working, the filament power is regulated according to average anode current. This paper designs an off-line single-switch zero-voltage-switching inverter as the filament power supply. As the transistor switch of the inverter is turned on, the input voltage drives the cathode filament via a step-down isolation transformer. As the transistor switch is turned off, the transformer resonant voltage drives the filament. The output power is dominated by the switch conduction time. The resonant parameters determine the switch off time.

This work is supported by the MOST of Taiwan; the contract number is MOST 107-2637-E-131-003.

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